

The effect of isoscalar-isovector coupling in infinite nuclear matter

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Working on the framework of Relativistic Mean Field theory, we exposed the effect of nonlinear isoscalar-isovector coupling on G2 parameter set on the density dependence of nuclear symmetry energy in infinite nuclear matter. The observables like symmetric energy and few related coefficients are studied systematically. We presented the results of stiff symmetry energy at sub-saturation densities and a soft variation at normal densities. Correlation between the symmetric energy and the isoscalar-isovector coupling parameter fully demonstrated for wide range of density. The work further extended to the octet system and showed the effect of coupling over the equation of state.

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I. INTRODUCTION

The best possible and well defined theoretical laboratory to studied for many body system is infinite nuclear matter at certain conditions. To review the status of microscopic studies of nuclear matter and neutron rich matter, that already reaches to the destination by mean-field models (relativistic and non-relativistic) [1–6] and although by some other microscopic, methods such as Brueckner-Hartree-Fock (BHF) [7–10] and Dirac-Brueckner-Hartree-Fock (DBHF) [11–15]. The isospin and density dependence of the nuclear symmetry energy, E_{sym} , is the current interest because of its implications not only in nuclear physics but also in astrophysics. The softening in equation of state (EOS) of nuclear matter likely lead to an exciting problem in astro-nuclear physics from few decades to till date. Mean while the novel phenomena like formation of superheavy nuclei in astrophysical system is the front learning about the island of stability for that regions.

The prediction of compressional modulus (K_0) in symmetric nuclear matter for nonrelativistic and relativistic models differ by about 25 in magnitude. In other word, the nonrelativistic models predict $K_0 \sim 220\text{--}235$ MeV [4–6], relativistic models argue for $K_0 \sim 250\text{--}280$ MeV [1–3]. The density dependence of the symmetry energy E_{sym} is not exposed fully at present, which may be responsible for the above discrepancy [16]. Note that the E_{sym} is the difference between the energy of pure neutron matter and normal nuclear matter, which is a tremendous approximation [3, 17, 18]. The large value of K_0 shows appropriately stiff symmetric energy i.e. rises rapidly with baryon density [16, 17]. This make a passing reference between K_0 and E_{sym} , which is a function of density and isospin component of scalar and vector mesons. Thus, it is enlightening to explore the effects of E_{sym} in the RMF model, which was limited to a narrow range from the analysis of skin data, on the composition and structure of hot proton neutron and cold neutron star matter that hold a large density range [19–23].

In this work, we introduce an extra term on the top of the G_1 or G_2 forces to the Lagrangian, which is the combined effect of scalar and vector field with isoscalar- isovector coupling constant (Λ_v). The newly added term affect to the sym-

metric energy and the slope parameter L_{sym} . The problem is the stiffness of the E_{sym} with respect to baryonic density followed a softer path with respect to density along with this extra added coupling. Without this additional coupling constants one may not be overcome the hindrance in the nuclear matter. The most important prospectus of the coupling constant is that the various other properties and nuclear matter observables are not seriously compromising in their magnitude.

II. THE RELATIVISTIC MEAN FIELD (RMF) THEORY

The complete description of EOS can be implemented by using quantum chromodynamics (QCD), which is the fundamental theory of the strong interaction. The concept of effective field theory (EFT) is applicable to low energy QCD [24] and is called as quantum hadrodynamics (QHD) [25–29]. The description of normal nuclear matter properties [25, 30] and finite nuclei [31–34] are successfully explained by QHD. The exchange of mesons like iso-scalar scalar σ , iso-scalar vector ω , iso-vector vector ρ and the pseudoscalar π are responsible for nucleon interaction and increases the degree of freedom. The nucleons are considered as Dirac spinors moving in classical meson fields. In mean-field approximation, the contribution of the π meson is zero, because of its pseudospin nature. The extension of the standard relativistic mean-field (RMF) theory with addition of non-linear scalar-vector and vector-vector self interaction is called as E-RMF proposed by *Furnstahl, Serot, and Tang* [35, 36] (see also *Miller and Serot* [37]). The coupling constants are self-consistent with the underlying symmetries of QCD by non-renormalization. The concept of naturalness can focus by expanding the non-linear Lagrangian and organized it in increasing power of field and their derivative according to their dimensional analysis [38, 39] and truncated at a given level of accuracy [40–42]. Here we have taken the fourth order of expanded field and in addition to this we have taken the special interaction between the ω and ρ with a new coupling constant Λ_v to get good result in the study of E_{sym} and other co-efficients. The modified Lagrangian is given as [43, 44]:

$$\begin{aligned} \mathcal{L} = & \sum_B \bar{\psi} (\{i\gamma^\mu D_\mu - m_B + g_\sigma \sigma\} \psi \\ & + \frac{1}{2} \partial_\mu \sigma \partial_\mu \sigma - m_\sigma^2 \sigma^2 \left(\frac{1}{2} + \frac{\kappa_3}{3!} \frac{g_\sigma \sigma}{m_B} + \frac{\kappa_4}{4!} \frac{g_\sigma^2 \sigma^2}{m_B^2} \right) \\ & - \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu \left(1 + \eta_1 \frac{g_\sigma \sigma}{m_B} + \frac{\eta_2}{2} \frac{g_\sigma^2 \sigma^2}{m_B^2} \right) \\ & - \frac{1}{4} R_{\mu\nu}^a R^{\mu\nu a} + \frac{1}{2} m_\rho^2 \rho_\mu^a \rho^{a\mu} \left(1 + \eta_\rho \frac{g_\sigma \sigma}{m_B} \right) \\ & + \frac{1}{4!} \zeta_0 g_\omega^2 (\omega_\mu \omega^\mu)^2 + \Lambda_v (g_\rho^2 \rho_\mu^a \rho^{a\mu}) (g_\omega^2 \omega_\mu \omega^\mu) \end{aligned} \quad (1)$$

The co-variant derivative D_μ is defined as

$$D_\mu = \partial_\mu + i g_\omega \omega_\mu + i g_\rho I_3 \tau^a \rho_\mu^a, \quad (2)$$

where $R_{\mu\nu}^a$ and $\Omega_{\mu\nu}$ are field tensors and defined as follow

$$R_{\mu\nu}^a = \partial_\mu \rho_\nu^a - \partial_\nu \rho_\mu^a + g_\rho \epsilon_{abc} \rho_\mu^b \rho_\nu^c, \quad (3)$$

where, $E_B^* = \sqrt{(k^2 + m_B^*)}$ is the effective energy, k_B is the Fermi momentum, $m_B^* = m_B - g_\sigma \sigma$ stands for effective mass (which solved self-consistently), I_3 is the third component of isospin projection and B stands for baryon. The energy and pressure density depends on the effective mass m_B^* of the system, so first need to solve this self-consistent Eqn. (4). As the self-consistent field obtained, the expression for energy and pressure density are given as;

$$\begin{aligned} \mathcal{E} = & \sum_B \frac{2}{(2\pi)^3} \int_0^{k_B} d^3 k E_B^*(k) + \frac{1}{8} \zeta_0 g_\omega^2 \omega_0^4 \\ & + m_\sigma^2 \sigma_0^2 \left(\frac{1}{2} + \frac{\kappa_3}{3!} \frac{g_\sigma \sigma}{m_B} + \frac{\kappa_4}{4!} \frac{g_\sigma^2 \sigma^2}{m_B^2} \right) \\ & + \frac{1}{2} m_\omega^2 \omega_0^2 \left(1 + \eta_1 \frac{g_\sigma \sigma}{m_B} + \frac{\eta_2}{2} \frac{g_\sigma^2 \sigma^2}{m_B^2} \right) \\ & + \frac{1}{2} m_\rho^2 \rho_0^2 \left(1 + \eta_\rho \frac{g_\sigma \sigma}{m_B} \right) \\ & + 3 \Lambda_v \Omega_0^2 R_0^2 + \sum_l \varepsilon_l, \end{aligned} \quad (10)$$

and

$$\Omega_{\mu\nu} = \partial_\mu \omega_\nu - \partial_\nu \omega_\mu, \quad (4)$$

Here m_B is the mass of nucleon, and m_σ , m_ω and m_ρ are the masses of corresponding mesons. The equation of motion derived from the Lagrangian using mean-field approximation self-consistently. The field equation for different meson fields are given by

$$\begin{aligned} m_\sigma^2 \left(\sigma_0 + \frac{g_\sigma \kappa_3 \sigma_0^2}{2m_B^2} + \frac{g_\sigma^2 \kappa_4 \sigma_0^3}{6m_B^2} \right) \\ - \frac{1}{2} m_\omega^2 \left(\eta_1 \frac{g_\sigma}{m_B} + \eta_2 \frac{g_\sigma^2 \sigma_0}{m_B^2} \right) \omega_0^2 \\ - \frac{1}{2} m_\rho^2 \eta_\rho \frac{g_\sigma \rho_0^2}{m_B} = \sum g_\sigma \rho_s, \end{aligned} \quad (5)$$

$$\begin{aligned} m_\omega^2 \left(1 + \eta_1 \frac{g_\sigma \sigma_0}{m_B} + \eta_2 \frac{g_\sigma^2 \sigma_0^2}{m_B^2} \right) \omega_0 \\ - \frac{1}{6} \zeta_0 g_\omega^2 \omega_0^3 + 2 \Lambda_v g_{rho}^2 \rho_0^2 g_\omega^2 \omega_0 = \sum g_\omega \rho_B, \end{aligned} \quad (6)$$

$$\left[m_\rho^2 \left(1 + \eta_\rho \frac{g_\sigma \sigma_0}{m_B} \right) + 2 \Lambda_v g_{rho}^2 g_\omega^2 \omega_0^2 \right] \rho_0 = \frac{1}{2} \sum g_\rho I_3 \rho \quad (7)$$

The scalar (ρ_s) and vector (ρ) densities are given as follow:

$$\rho_s = \frac{1}{\pi^2} \int_0^{k_B} \frac{k^2 dk}{E_B^*} \quad (8)$$

and

$$\rho_B = \frac{1}{\pi^2} \int_0^{k_B} k^2 dk \quad (9)$$

$$\begin{aligned} \mathcal{P} = & \sum_B \frac{2}{3(2\pi)^3} \int_0^{k_B} d^3 k E_B^*(k) + \frac{1}{8} \zeta_0 g_\omega^2 \omega_0^4 \\ & + m_\sigma^2 \sigma_0^2 \left(\frac{1}{2} + \frac{\kappa_3}{3!} \frac{g_\sigma \sigma}{m_B} + \frac{\kappa_4}{4!} \frac{g_\sigma^2 \sigma^2}{m_B^2} \right) \\ & + \frac{1}{2} m_\omega^2 \omega_0^2 \left(1 + \eta_1 \frac{g_\sigma \sigma}{m_B} + \frac{\eta_2}{2} \frac{g_\sigma^2 \sigma^2}{m_B^2} \right) \\ & + \frac{1}{2} m_\rho^2 \rho_0^2 \left(1 + \eta_\rho \frac{g_\sigma \sigma}{m_B} \right) \\ & + \Lambda_v \Omega_0^2 R_0^2 + \sum_l P_l, \end{aligned} \quad (11)$$

where P_l and ε_l are lepton's pressure and energy density, respectively (and here we have taken electron only). Note that the following quantities are redefined as, $\Omega_0 \equiv g_\omega \omega_0$ and $R_0 \equiv g_\rho \rho_0$. At zero temperature, the pressure of the system obtained from energy density and its first derivative, which is thermodynamically consistent in mean-field (MF) approximation.

III. CALCULATIONS AND DISCUSSION

We have solved the mean field equations self-consistently and parametrized the coupling constant by renormalization. We have made attention to G2 [35, 36] parametrization sets of the E-RMF model to reproduce the fundamental properties of nuclear matter. In this parameter set ζ coupling constant is determined by self-interactions of ω -meson, which is responsible for the softening the EOS at high density and reproduced the maximum mass of neutron star. For softening the symmetric energy of the system at nuclear matter density,

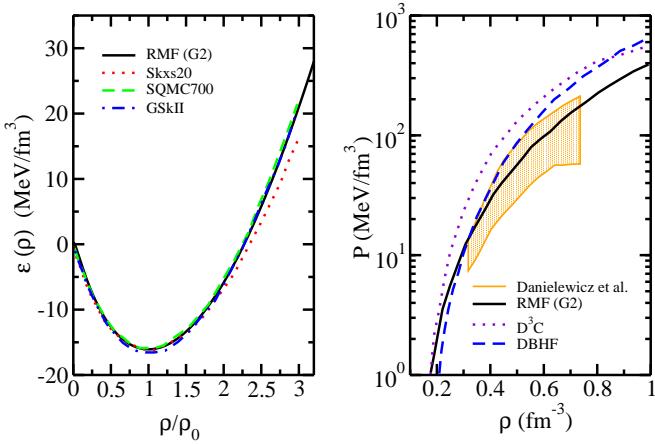


FIG. 1: The EOS or energy per nucleon and pressure density as a function of density from RMF compare with the other predictions.

we have added an extra coupling between iso-scalar and iso-vector field with coupling constant Λ_v . This made make possible to soften the E_{sym} and reproduced saturation properties for the force parameter at $\Lambda_v \sim 0.00$. In other word, Λ_v is simply a self interaction constant between isoscalar and isovector field, which does not change any other nuclear matter properties like energy density, pressure density and so many considerably.

A. Energy and Pressure density

Infinite nuclear matter is important for the investigation of physical quantities relevant to heavy nuclei and compact objects like neutron star. At saturation density ρ_0 , the binding energy per particle ϵ as a function of density is well established quantities consistent to thermodynamic relation at certain temperature. These translational invariance of the system facilitates theoretical calculations of these observables. The prediction of E-RMF ($G2+\Lambda_v$) for EOS of nuclear matter (NM) (solid black line) compared with DBHF [45] and other theoretical predictions [46] are in Fig. 1 (a). From the figure, it is clear that all the prediction well matches with respect to density. Fig. 1(b), shows the comparison of pressure density with experimental data [47] and theoretical predictions [45, 48, 49]. The obtained results for energy density (see Fig. 1) from Ref. [21] is exactly same to the present calculation. This implies that the EOS of SNM is not affected by the new coupling with Λ_v . Being insensitive to this observable, the aim of this present work to pursue the systematic variation of E_{sym} by employing the isoscalar-isovector coupling parameter Λ_v , which discussed below.

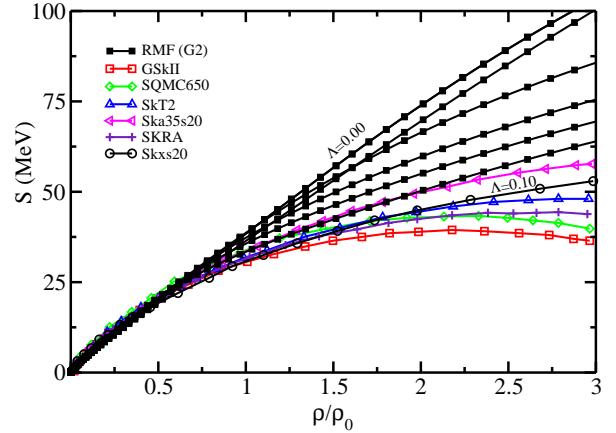


FIG. 2: The results of symmetric energy E_{sym} is a function of density ρ from RMF for different values of Λ_v compare with others predictions.

B. Symmetry Energy

The symmetric energy E_{sym} , which is important in infinite nuclear matter and finite nuclei because of isospin dependent in the interaction. The expression of this key observable is defined as

$$E_{sym}(\rho) = \frac{k_F^2}{6E_F^*} + \frac{g_\rho^2 \rho_B}{2} \left(\frac{1}{4m_\rho^{*2}} + \frac{\Lambda_v g_\rho^2 g_\omega^2 \omega^2}{m_\rho^{*2}} \right), \quad (12)$$

where the effective energy $E_F^* = \sqrt{(k_F^2 + m^{*2})}$, k_F is the Fermi momentum and the effective mass $m^* = m - g_S \phi$. The effective mass of the ρ -meson modified because of non-linear coupling ($\rho - \omega$) interaction and is given by

$$m_\rho^{*2} = \left(1 + \eta_{rho} \frac{g_\sigma \sigma}{m_B} \right) m_\rho^2 + 2g_\rho^2 (\Lambda_v g_v^2 \omega_0^2), \quad (13)$$

where m_ρ is the mass of the ρ -meson. It is noted that the non-linear isoscalar-isovector coupling (Λ_v) modified the density dependent of E_{sym} without affecting the saturation properties of the SNM (as shown in Fig. 1 (a)). In this model the symmetric nuclear matter saturates at a Fermi momentum of $k_F = 1.00 \text{ fm}^{-1}$ with $E \sim 16.0 \text{ MeV}$ and an compressibility of $K_0 = 215 \text{ MeV}$ [21]. We can expand the density dependence of the E_{sym} with the help of these 3-co-efficients as a function of $\mathcal{Y} = (\rho - \rho_0) / 3\rho_0$ as follow:

$$\begin{aligned} \mathcal{S}(\rho) = & E_{sym}(\rho_0) + L_{sym} \mathcal{Y} + \frac{1}{2} K_{sym} \mathcal{Y}^2 \\ & + \frac{1}{6} Q_{sym} \mathcal{Y}^3 + O(\mathcal{Y}^4), \end{aligned} \quad (14)$$

where L_{sym} and K_{sym} are defined as

$$L_{sym} = 3\rho_0 \left(\frac{\partial E_{sym}}{\partial \rho} \right)_{\rho=\rho_0}, \quad (15)$$

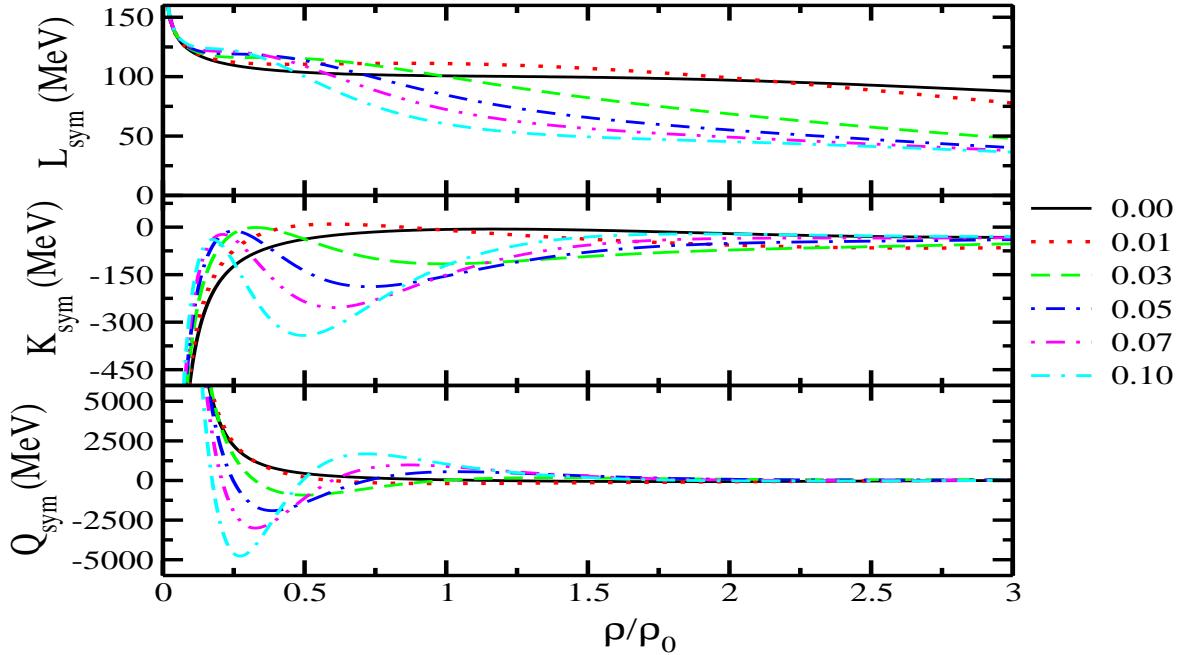


FIG. 3: The slope parameter L_{sym} along with K_{sym} and Q_{sym} for different value of Λ_v , with G2 parameter set.

$$K_{sym} = 3 \left(\rho_0 \frac{\partial L_{sym}}{\partial \rho} - L_{sym} \right)_{\rho=\rho_0}. \quad (16)$$

and

$$Q_{sym} = 27 \rho_0^3 \left(\frac{\partial^3 E_{sym}}{\partial \rho^3} \right)_{\rho=\rho_0}. \quad (17)$$

Here L_{sym} is called the slope parameter, which is nothing but the slope of E_{sym} . The quantity K_{sym} represents as the curvature of E_{sym} with respect to density. In SNM, the value of K_{sym} at saturation density (ρ_0) entitled by K_0 and called as symmetry incompressibility. Finally, Q_{sym} is the third derivative of the symmetric energy, which called as skewness coefficient [46]. As shown in the Fig. 2, we plotted the E_{sym} for RMF (G2+ Λ_v) compare with other theoretical predictions [46]. The softness of the E_{sym} with the coupling Λ_v clearly observed from the figure. Further, the rate of change of E_{sym} i.e. the slope parameter L_{sym} for different values of Λ_v is displayed in Fig. 3. Both the figure shows similar behaviour with the coupling constant and with baryonic density.

There are many references which trying to fix the constraint on symmetric energy and L_{sym} [46, 50–55]. Finally we comes to the constraint on E_{sym} and L_{sym} at saturation density (ρ_0) and the obtained results are compared with with experimental data at HIC [56], PDR [57, 58], IAS [59] and FRDM [60] in Fig. 4 and Table 2. From the figure, it is clear that for $\Lambda_v=0.00$, our observables are not within the experimental constraint region. For the values of $\Lambda_v=0.01-0.04$,

TABLE I: The results for slope co-efficient S_{sym} , K_{sym} for different value of Λ_v are in list.

Λ_v	E_{sym}	L_{sym}
0.00	36.40	105.67
0.01	38.41	111.69
0.02	38.87	113.09
0.03	38.62	112.32
0.04	38.02	110.52
0.05	37.27	108.27
0.06	36.46	105.85
0.07	35.64	103.42
0.08	34.86	101.05
0.09	34.11	98.79
0.10	33.39	96.65

these comes within limit and again go out with further increasing the value of $\Lambda_v = 0.05$. Hence, from this behavior, we may suggest the limiting value of Λ_v is 0.01-0.04, which can improved the observables with G2 force parameter.

IV. HADRON MATTER

The density at the interior core of the neutron star is very high $\sim 7-8\rho_0$ so other hadronic part [43, 44, 61–63] also come to play role on the equation of state. In this section,

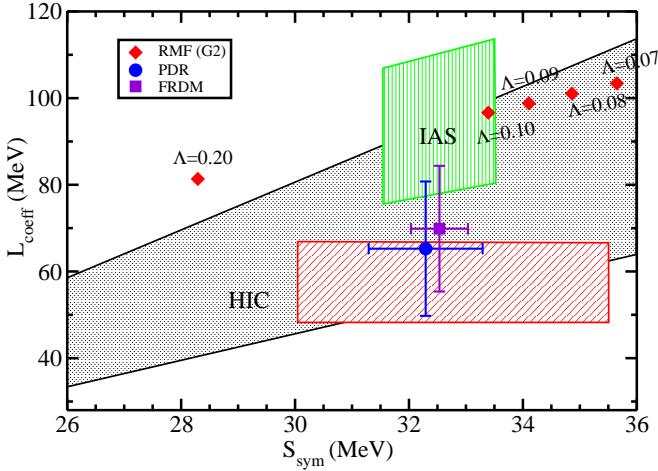


FIG. 4: The Symmetric energy S_0 with respect to L_{sym} at saturation density compare with flow data.

we extended our calculations to the system with octet ($n, p, \Lambda^0, \Sigma^{0,\pm 1}, \Xi^{-0}$). The purpose of this work to show the effect of isoscalar-isovector interaction with coupling constant Λ_v on the equation of state of the nuclear system with full octet. It is worthy to mention that, many authors [21, 41, 64] have been used $G1$ and $G2$ forces for the description of finite and infinite nuclear matter properties. Moreover, the non-linear isoscalar-isovector interaction in mean field was discussed in Refs. [65, 66], here we suppose to check this interaction with $G2$ along with Λ_v , keeping all saturation constraints. We assume that all the hyperon in the octet system have the same coupling ratio with mesons. This is given by a constant factor which is equal to $\sqrt{2/3}$ [64], the density of the system $\sim 7\text{-}8 \rho_0$ and assuming that the stars are the composite system of baryons. The electron e^{-1} and muon μ are for maintaining the charge neutrality and $\beta-$ equilibrium condition [64, 67] under the week interactions:

$$B_1 \rightarrow B_2 + l + \bar{\nu}_l; \quad B_2 + l \rightarrow B_1 + \nu_l,$$

where $B_1, B_2, l, \nu, \bar{\nu}$ are baryons, leptons, neutrino and antineutrino respectively. In case of octet system, we are dealing with neutron star (real system), which became unstable with small change in the system. The equation of state obtained for nuclear matter with nucleon and hadrons for $G2 + \Lambda_v$ is shown in Fig 5.

We plotted pressure(fm^{-4}) density with respect to energy density (fm^{-4}) in Fig. 5. Here we can see that at high density octet EOS becomes stiffer by increasing the Λ_v value but it behave different way at low density region. Equation of state for nucleon case (EOS-NP) becomes softer with increasing Λ_v value over whole the density region. We compare our results with the empirical EOS obtained by Steiner *et al.* [68]. Along with the octet we also plotted nucleon EOS at different Λ_v values. Nucleon EOS well matches with the data of Steiner *et al.* [68] at low and high density but the octet EOS is devi-

ates at high density. The octet EOS compares well with the

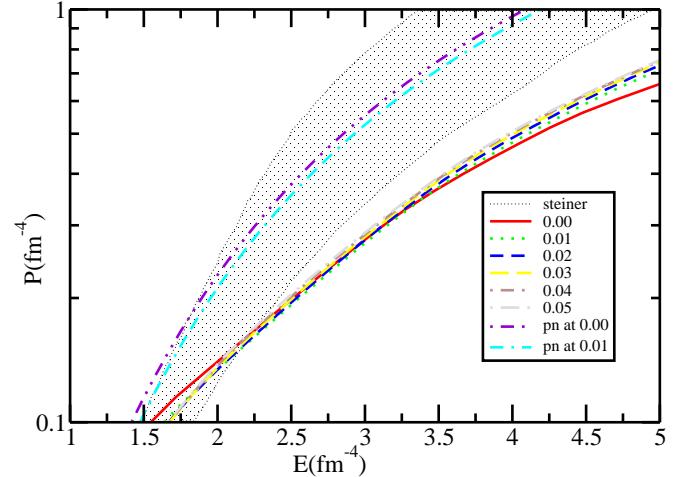


FIG. 5: Equation of states(EOS) with the E-RMF model for octet and NP at different Λ_v value.

experimental data only at intermediate density. An impressive observation is noted from this plot, the influence of Λ_v with $G2$ parameter set is not much because in $G2$ already taken the four power of iso-scalar meson so it is more effective then iso-scalar and iso-vector interaction.

V. SUMMARY AND CONCLUSIONS

After analyzing the result whatever got from this calculation, we forcefully say that Λ_v i.e. the coupling of isoscalar and isovector interaction is very crucial to softening the equation of state of nucleonics and octet system. One can easily observed that the E_{sym} curve with respect to ρ/ρ_0 became softer with varying Λ_v but after $\Lambda_v \sim 0.05$ it became constant. Same result we got in L_{sym} which is the slope of symmetric energy. For symmetric nuclear matter isoscalar and isovector interaction do not effect the E/A and pressure density curve with respect to ρ/ρ_0 . We got identical curve as in Ref. [21] for $G1$ and $G2$ sets with experimental data [47]. In case of hadron, the extra term containing Λ_v plays an important role in softening the equation of state. Hence, we may say that this interaction term is crucial for both cases (nucleon and hadron). Further, we extend our calculation to the neutron star properties in near future.

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